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ANALYZING POWER MEASUREMENTS FOR THE EXCITATION OF STATES IN  $^{28}\text{Si}$  and  $^{24}\text{Mg}$  BY INELASTIC SCATTERING OF POLARIZED PROTONS

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Analyzing powers  $A_y(\theta)$  for the excitation of states in  $^{28}\text{Si}$  and  $^{24}\text{Mg}$  with excitation energies up to 16 MeV have been measured with a 135-MeV polarized proton beam. The scattered protons were detected with the QDDM magnetic spectrograph at angles between  $25^\circ$  and  $65^\circ$  with an overall resolution of about 70 keV. Results for the  $6^-$ ,  $T=1$  (14.35 MeV),  $6^-$ ,  $T=0$  (11.58 MeV), and  $5^-$ ,  $T=0$  (9.70 MeV) states in  $^{28}\text{Si}$ ,<sup>1</sup> whose predominant configurations are all  $(d_{5/2})^{-1}(f_{7/2})$ , are shown in Fig. 1, where they are compared with the results of DWIA calculations using the  $t$ -matrix effective interaction derived by Love from the free nucleon-nucleon scattering data.<sup>2</sup> The cross section for the  $6^-$ ,  $T=1$ , state is due mainly to the tensor direct term in the interaction, while that for the  $6^-$ ,  $T=0$ , state is due mainly to tensor and spin-orbit exchange terms, and that for the  $5^-$ ,  $T=0$ , state is due mainly to spin-orbit and central interaction terms. The  $A_y(\theta)$  results for the  $6^-$  states are sensitive to interference both between the central and spin-orbit parts and between the spin-orbit and tensor parts of

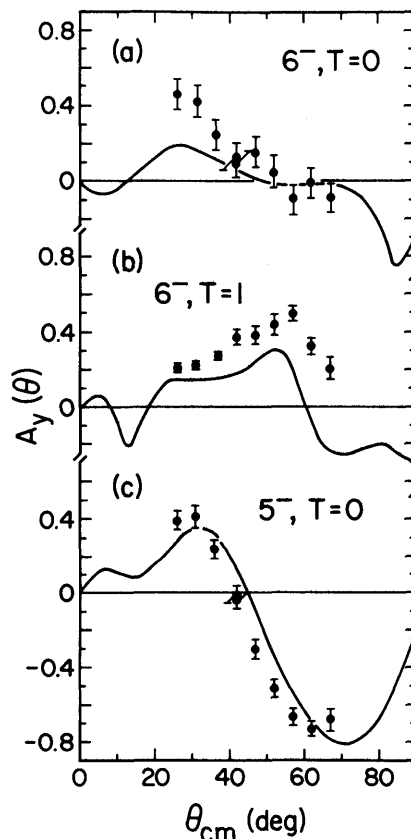


Figure 1. Analyzing powers,  $A_y(\theta)$ , for the 135-MeV  $(\bar{p}, p')$  excitation of (a) the  $6^-$ ,  $T=0$ , state at 11.58 MeV, (b) the  $6^-$ ,  $T=1$ , state at 14.35 MeV, and (c) the  $5^-$ ,  $T=0$ , state at 9.70 MeV. The experimental data are compared with results of DWIA calculations using the Love  $t$ -matrix.

the t-matrix. For the  $5^-$  state, on the other hand, the calculated result for  $A_y$  is predominantly sensitive to central-spin-orbit interference, and the change of sign of  $A_y$  near  $40^\circ$  is rather well correlated with the change of sign in the central part of the t-matrix at the corresponding momentum transfer, as given by the Love interaction.

The differential cross sections and analyzing powers for the lower-lying states of  $^{28}\text{Si}$  and  $^{24}\text{Mg}$  have also been measured. In  $^{24}\text{Mg}$ , for example, the excitation of states in the  $K=0$  and  $K=2$  bands can be compared to calculations using the Chalk River projected-Hartree-Fock wave functions. The remarkable agreement between theory and experiment for inelastic

electron scattering<sup>3</sup> allows a detailed study of the proton scattering mechanism. The totally anomalous shape of the electromagnetic form factor for the  $4_1^+(K=0)$  state in  $^{24}\text{Mg}$  is well reproduced in the differential cross section for (p,p').

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SPIN-ORBIT EFFECTS IN THE EXCITATION OF PROTON AND NEUTRON STATES  
IN THE (p,p') REACTION AT 160 MeV, 120 MeV, AND 95 MeV

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Large differences in the shapes of measured differential cross sections were found earlier<sup>1</sup> for the excitations of the  $4_1^+$  proton state in  $^{90}\text{Zr}$  and the  $4_1^+$  neutron state in  $^{92}\text{Zr}$ . To obtain a good fit to the data for this proton state in  $^{90}\text{Zr}$  with purely collective calculations, an enhanced spin-orbit contribution ( $\beta_4^{80}/\beta_4=1.25$ ) was required, but no satisfactory fits were found for this neutron state in  $^{92}\text{Zr}$ . Collective fits to the data for the  $2_1^+$ ,  $4_1^+$ ,  $6_1^+$ ,  $8_1^+$  proton states in  $^{90}\text{Zr}$  showed<sup>2</sup> the increasing dominance of the spin-orbit contribution as the multipolarity increased. Recent calculations show the cross sections for the  $2_1^+$ ,  $4_1^+$ ,  $6_1^+$ , and  $8_1^+$  states in  $^{90}\text{Zr}$  to be underpredicted in the DWIA by factors of 30, 10, 3, and 2 respectively when only the  $(g_{9/2})^2$  valence terms are included for

the central, spin-orbit and tensor amplitudes,<sup>3</sup> suggesting the need for core polarization amplitudes similar in magnitude to those required at lower energies.<sup>4,5</sup> These DWIA calculations showed the relative importance of the spin-orbit part of the t-matrix increased in this sequence as the multipolarity increased.<sup>3</sup> The dominance of the spin-orbit contributions for the  $8_1^+$  state in  $^{90}\text{Zr}$  is shown in the DWIA and collective calculations of Figures 1(a) and 1(b) respectively.

Large spin-orbit effects at this energy ( $E_p=160$  MeV) clearly suggested the need for (p,p') asymmetry measurements. Analyzing power data have been obtained at 14 angles from  $16^\circ$  to  $44^\circ$  for the  $2_1^+$ ,  $4_1^+$ ,  $5_1^-$ ,  $3_1^-$ , and  $2_2^+$  states in  $^{90}\text{Zr}$  and for the  $2_1^+$ ,  $4_1^+$ ,  $3_1^-$ ,  $2_2^+$ ,  $2_3^+$ , and  $5_1^-$  states in  $^{92}\text{Zr}$ , and at 8 angles from  $26^\circ$