

of 100 and 39 m, respectively. For the ^{208}Pb target, measurements at only a few angles have been obtained. In the present work, spectra have been measured at $E_p=160$ MeV in order to further investigate the nature of these excitations and of the continuum.

Figure 1 shows a sample A_y spectrum for the $^{90}\text{Zr}(p,n)^{90}\text{Nb}$ reaction at $\theta=4.2^\circ$. The indicated excitation regions associated with each previously identified state have distinct values of A_y . At these

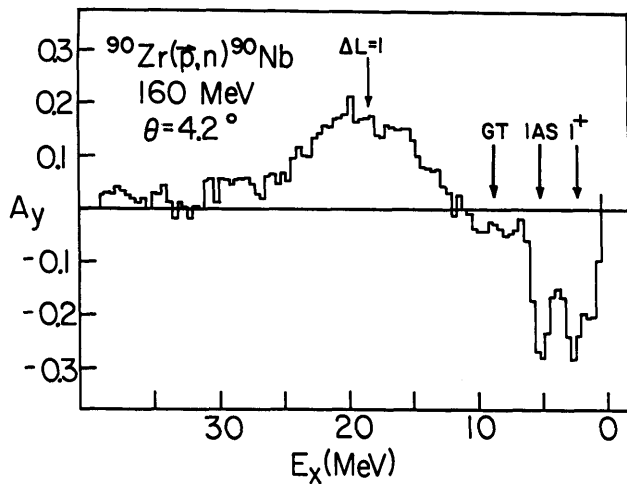


Figure 1. Analyzing power for the $^{90}\text{Zr}(p,n)^{90}\text{Nb}$ reaction at 4.2° . The uncertainty in A_y is about ± 0.02 .

forward angles the continuum, does not appear to have a significantly non-zero A_y . Preliminary extraction of the A_y data for the lower 1^+ state differ from that for

the giant GT, possibly reflecting the expected difference in the dominant particle-hole configurations. Alternately the A_y of the broad $\Delta L=1$ resonance do not appear to exhibit a strong excitation energy dependent variation consistent with discrete concentrations of 0^- , 1^- , and 2^- strength.⁵ The back angle A_y spectra are dominated by the continuum. Data have been obtained over an excitation range of 0 to 60 MeV. As the angle increases from 15° to 48° , the A_y become increasingly more positive. The largest A_y at each angle are associated with the lowest excitation energies, with the maximum A_y reaching a value of about +0.5 at $\theta=48^\circ$.

Further reduction of the data and DWIA calculations are in progress.

- 1) C. Gaarde, J. Rapaport, T.N. Taddeucci, C.D. Goodman, C.C. Foster, D.E. Bainum, C.A. Goulding, M.B. Greenfield, D.J. Horen, and E. Sugarbaker, Nucl. Phys. A369, 258 (1981).
- 2) D.E. Bainum, J. Rapaport, C.D. Goodman, D.J. Horen, C.C. Foster, M.B. Greenfield, and C.A. Goulding, Phys. Rev. Lett. 44, 1751 (1980).
- 3) D.J. Horen, C.D. Goodman, C.C. Foster, C.A. Goulding, M.B. Greenfield, J. Rapaport, D.E. Bainum, E. Sugarbaker, T.G. Masterson, F. Petrovich, and W.G. Love, Phys. Lett. 95B, 27 (1980).
- 4) D.J. Horen, C.D. Goodman, D.E. Bainum, C.C. Foster, C.C. Gaarde, C.A. Goulding, M.B. Greenfield, J. Rapaport, T.N. Taddeucci, E. Sugarbaker, T. Masterson, S. Austin, A. Galonsky, and W. Sterrenburg, Phys. Lett. 99B, 383 (1981).
- 5) F. Osterfeld, S. Krewald, H. Dermawan, and J. Speth, Phys. Lett. 105B, 257 (1981).

STUDY OF THE $(d,^2\text{He})$ REACTION AT $E_d = 99$ MeV

K.B. Beard, J. Kasagi, E. Kashy and B.H. Wildenthal
Cyclotron Laboratory, Michigan State University, East Lansing, Michigan 48824

D.L. Friesel and H. Nann
Indiana University Cyclotron Facility, Bloomington, Indiana 47405

R.E. Warner
Michigan State University and Oberlin College, Oberlin, Ohio 44074

The ability to measure the strengths of $\Delta L=0$, $\Delta J=\Delta S=\Delta T=|\Delta T_z|=1$ transitions in the T_z direction of

increasing neutron excess would be of great general utility in probing nuclear structure and would have

special interest in answering problems such as that of determining the spectrum for electron capture on ^{56}Fe in the astrophysical environment leading to the formation of neutron stars.¹ The reaction $(d, ^2\text{He})$, in which the reaction product is the "di-proton" singlet state with $T=1$ and $S=0$, is in principle ideally suitable as a probe for such transitions. It provides a charge exchange mode induced by a charged, non-radioactive beam; moreover, due to the quantum numbers of the initial and final two-nucleon states, it automatically eliminates the non-spin-flip component which is admixed with the spin-flip mode in the reactions initiated with half-integer-spin projectiles. The internal structure of the deuteron and di-proton systems should be much simpler than those of heavier complex projectiles such as ^{12}C and ^{14}N which also satisfy these criteria.

The demonstrated feasibility² of detecting the unbound ^2He system as a reaction product in the $(\alpha, ^2\text{He})$ reaction immediately suggests exploring the $(d, ^2\text{He})$ reaction as a spin-flip, charge-exchange probe. Such a study was carried out at $E_d = 55$ MeV on ^6Li , ^{10}B and ^{12}C .³ The results of this study were positive in that the observed cross sections were consistent with a direct, one-step charge-exchange mechanism. That is, the shapes and magnitudes of the measured angular distribution were reasonably well matched by a combination of DWBA reaction-mechanism calculations and shell-model structure predictions. However, the angular distributions observed were not strikingly characteristic of particular angular momentum transfers and, in particular, the $\Delta L=0$ transitions were neither noticeably enhanced in magnitude nor easily identifiable by virtue of their shape. The results of the 55 MeV study thus suggest that the $(d, ^2\text{He})$ reaction leading to known final states can yield valuable

spectroscopic information about these states, but they do not demonstrate that this reaction can be successfully used as a probe of $\Delta S=1$, $\Delta T=1$ strength in a region of unknown structural features. To be successful in this latter mode, the reaction must supply a characteristic signature such as a dominant cross section for $\Delta L=0$ relative to other ΔL transfers or a combination of large cross section and distinctive shape of the angular distribution.

Motivated by the hope that the $(d, ^2\text{He})$ reaction might be more selective in enhancing $\Delta L=0$ spin-flip transitions at a higher bombarding energy, we have repeated the study of the $^{12}\text{C}(d, ^2\text{He})^{12}\text{B}$ reaction at 99 MeV. This particular reaction is convenient for the usual reasons which include target fabrication and stability of ^{12}C , the ground state spin of $J^\pi = 0^+$, and, more importantly, because the ground state of ^{12}B has $J^\pi = 1^+$. Its structure is known (from inelastic electron scattering studies⁴ connecting the ground state of ^{12}C to the $(J^\pi, T) = (1^+, 1)$ isobaric analog of the ^{12}B ground state) to have a large overlap with the ^{12}C target via the $\Delta S=1$, $\Delta T=1$ operator. Hence, if the $(d, ^2\text{He})$ reaction at 99 MeV is to be useful as probe for discovering $S=1$, $T=1$ strength, the ^{12}B ground state must dominate the final-state spectrum.

Our measurements employed a 99.2 MeV deuteron beam from the Indiana University Cyclotron Facility. Protons were detected in two solid-state detector telescopes (450 μ Silicon as ΔE , 15 mm Germanium as E) each subtending a solid angle of 1.14 msr. The detectors were mounted in the same vertical plane with their centers 2.33° above and below the horizontal plane.

Signals corresponding to the total energy deposited in each telescope (E_1 and E_2), the sum of the energy in the two telescopes ($E_1 + E_2$), the

particle identification from each telescope (PI₁ and PI₂) and the TAC-generated time difference between the E₁ and E₂ signals were recorded. The coincidence energy spectrum (E₁ + E₂) for ¹²C(d,²He)¹²B at 99.2 MeV and θ_{lab} = 20° is shown in Fig. 1. The energy

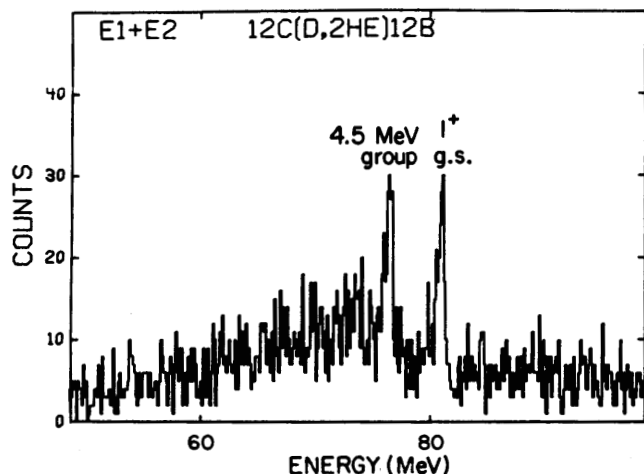


Figure 1. ²He energy spectrum from ¹²C(d,²He)¹²B reaction at E_d = 99 MeV and a laboratory angle of 20°.

resolution is 480 keV FWHM. The energy calibration indicates that the two strong peaks correspond to the J^π = 1⁺ ground state and a group of excited states (J^π = 4⁻ and 2⁻) at 4.5 MeV. In Fig. 2, the measured energy difference spectrum of the protons recorded in coincidence in the two detectors (|E₁-E₂|) is compared to the predictions of Watson-Migdal theory⁵ for the di-proton spectrum in our geometry. We take the correspondence as confirmation that our measurements do involve formation of a di-proton as the exit "particle". In Fig. 3 we show the angular distribution between 15° and 30° for the two strong peaks in the spectrum of Fig. 1. Both groups exhibit differential cross sections which decrease exponentially with increasing angle. The ground state group shows a steeper slope than that for the 4.5 MeV group.

Our results show little qualitative difference

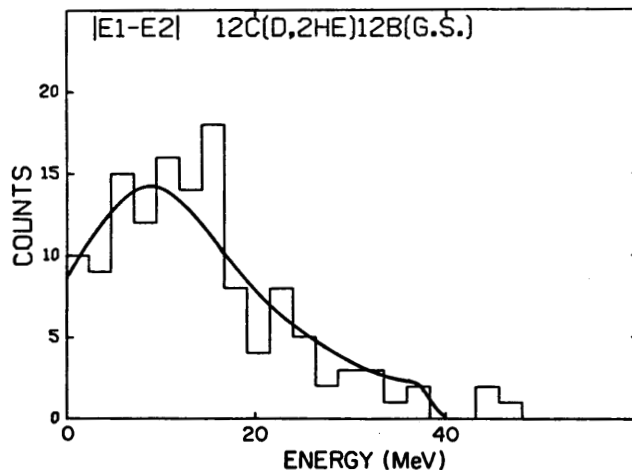


Figure 2. Energy difference spectrum for protons from the reaction ¹²C(d,²He)¹²B (g.s.) at E_d = 99 MeV. The solid line is the result of a Watson-Migdal final state interaction calculation that has been normalized to the data.

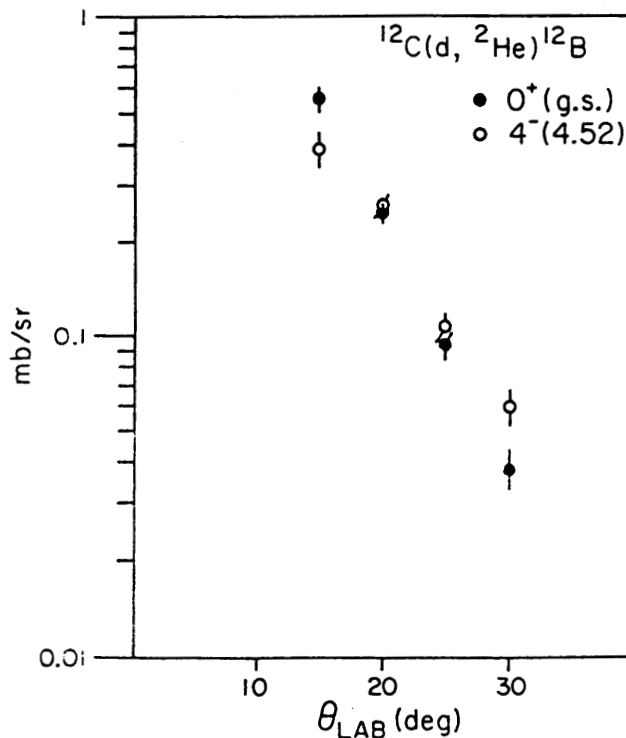


Figure 3. Cross section as a function of angle for ¹²C(d,²He)¹²B at E_d = 99 MeV.

from those obtained at lower energy. The angular distributions at 99 MeV are more strongly forward peaked than at 55 MeV, and the strength of the ground state relative to that of the first excited 2^+ state appears (within our statistics for the $J^\pi = 2^+$ state) larger at the higher energy. The former observation suggests of course, that measurements be extended to smaller angles in the hope that the cross sections of the $\Delta L=0$ transition would be, as is usually the case, further enhanced nearer 0° . Unfortunately, the deuteron beam carries an intrinsic impediment to such measurements in the form of the Coulomb dissociation of the deuteron. The cross section for this process increases with energy and with the approach to 0° .

The essential result of our measurements, however, is that the strength of the $J^\pi = 1^+$ ground state relative to that of the 4.5 MeV group is not

qualitatively different at 99 MeV than it was at 55 MeV. Hence, it does not appear feasible to use the ($d, {}^2\text{He}$) reaction at 99 MeV, any more than at 55 MeV, to probe a region of excitation in which level densities are high and specific spin assignments lacking as a means of identifying $\Delta L=0$, $\Delta S=1$, $\Delta T=1$ transitions and measuring their strengths.

- 1) H.A. Bethe, G.E. Brown, J. Applegate and J.M. Lattimer, Nucl. Phys. A324, 487 (1979).
- 2) R. Jahn, G.J. Wozniak, D.P. Stahel, and J. Cerny, Phys. Rev. Lett. 37, 812 (1976).
- 3) D.P. Stahel, R. Jahn, G.J. Wozniak and J. Cerny, Phys. Rev. C 20, 1680 (1979).
- 4) B.T. Chertok, C. Sheffield, J.W. Lightbody, Jr., S. Penner, and D. Blum, Phys. Rev. C 8, 23 (1973).
- 5) K.M. Watson, Phys. Rev. 88, 1163 (1952); A.B. Migdal, Sov. Phys. JETP 1, 2 (1955).
- 6) J.R. Wu, C.C. Chang, and H.D. Holmgren, Phys. Rev. C 19, 370 (1979); N. Matsuoka, M. Kondo, A. Shimizu, T. Saito, S. Nagamachi, H. Sakaguchi, A. Goto and F. Ohtai, Nucl. Phys. A345, 1 (1980).