

FERMI AND GAMOW-TELLER TRANSITIONS OBSERVED  
IN  $^{12,13,14}\text{C}(^3\text{He},t)^{12,13,14}\text{N}$  CHARGE EXCHANGE  
AT  $E(^3\text{He}) = 200 \text{ MeV}$

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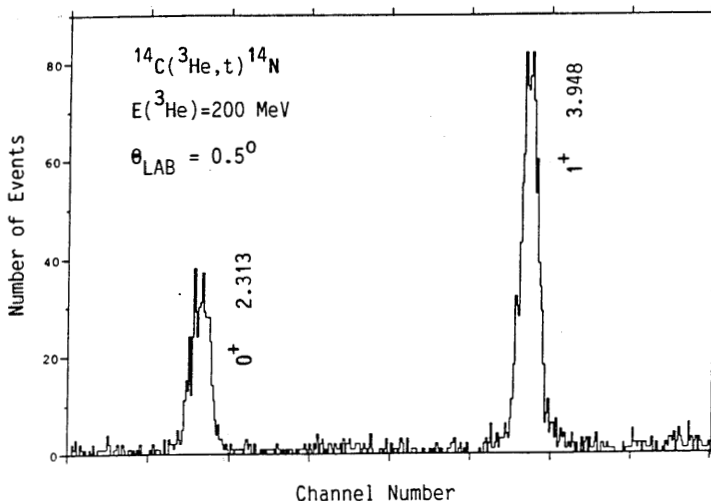
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The techniques established for performing measurements at  $\theta = 0^\circ$  with the K600 magnetic spectrometer have been used to determine  $(^3\text{He},t)$  cross sections for several low-excited states in  $^{12}\text{N}$ ,  $^{13}\text{N}$ ,  $^{14}\text{N}$ . Fig. 1 displays a spectrum for  $^{14}\text{C}(^3\text{He},t)^{14}\text{N}$ . The transitions to the  $0^+$ ,  $T=1$  (2.313 MeV) and  $1^+$ ,  $T=0$  (3.948 MeV) states represent pure Fermi and Gamow-Teller transitions, respectively.

Cross sections at  $\theta = 0^\circ$  are used extensively in  $(p,n)$  charge exchange to extract the ratio of spinflip to non-spinflip strength.<sup>1</sup> Using this procedure and our earlier data,<sup>2</sup> taken at  $E(^3\text{He}) = 200 \text{ MeV}$  for the transitions to the  $1/2^-$  (g.s.) and  $3/2^-$  (3.511 MeV) states in  $^{13}\text{N}$ , led to a ratio of spinflip to non-spinflip strength of about unity. Furthermore, it was shown from data for isobaric analog states from  $A=30$  to 208 that the non-spinflip strength  $V_\tau$  decreases for bombarding energies from 75 MeV to 200 MeV by a factor 0.6 in agreement with  $(p,n)$  data.



*Figure 1.* Triton energy spectrum from  $^{14}\text{C}(^3\text{He},t)^{14}\text{N}$  for the angular range  $= 0^\circ$  to  $1^\circ$  for transitions to the low-excited  $0^+$  and  $1^+$  states in  $^{14}\text{N}$ .

However, it is known<sup>3</sup> from results obtained at lower  $^3\text{He}$  bombarding energies that  $0^\circ$  cross sections are not sufficient to determine spinflip strength because of the presence of the tensor interaction  $V_{T\tau}$  in addition to the spinflip interaction  $V_{\sigma\tau}$ . It was, therefore, decided to measure complete angular distributions for several low-excited states with targets of  $^{12,13,14}\text{C}$ . Some of the preliminary results are displayed in Fig. 2. The non-spinflip Fermi transitions to the  $1/2^-$  (g.s.) state in  $^{13}\text{N}$  and the  $0^+$  (2.313 MeV) state in  $^{14}\text{N}$  show a more diffractive behavior at small angles compared to the spinflip Gamow-Teller transitions to the  $1^+$  (g.s.) state in  $^{12}\text{N}$ , the  $3/2^-$  (3.511 MeV) state in  $^{13}\text{N}$ , and the  $1^+$  (3.948 MeV) state in  $^{14}\text{N}$ . It should be noted, though, that the transition to the  $1/2^-$  (g.s.) state in  $^{13}\text{N}$  contains weak spinflip admixtures.

The theoretical interpretation of the data has not been completed. It will follow the earlier one-step DWBA analysis<sup>3</sup> performed for 65/90 MeV data. Microscopic wave functions will again be used, and the effective  $^3\text{He}$ -nucleon interaction will be parameterized with potentials of Yukawa shape and strengths  $V_\tau$ ,  $V_{\sigma\tau}$  and  $V_{T\tau}$ . It is expected that the tensor interaction will play a major role, and a comparison with the (p,n) charge-exchange reaction will become possible.

1. E. Sugarbaker, *et al.*, Phys. Rev. Lett. **65**, 551 (1990).
2. J. Jänecke, *et al.*, Nucl. Phys. **A526**, 1 (1991).
3. S. van der Werf, *et al.*, Nucl. Phys. **A496**, 305 (1989).

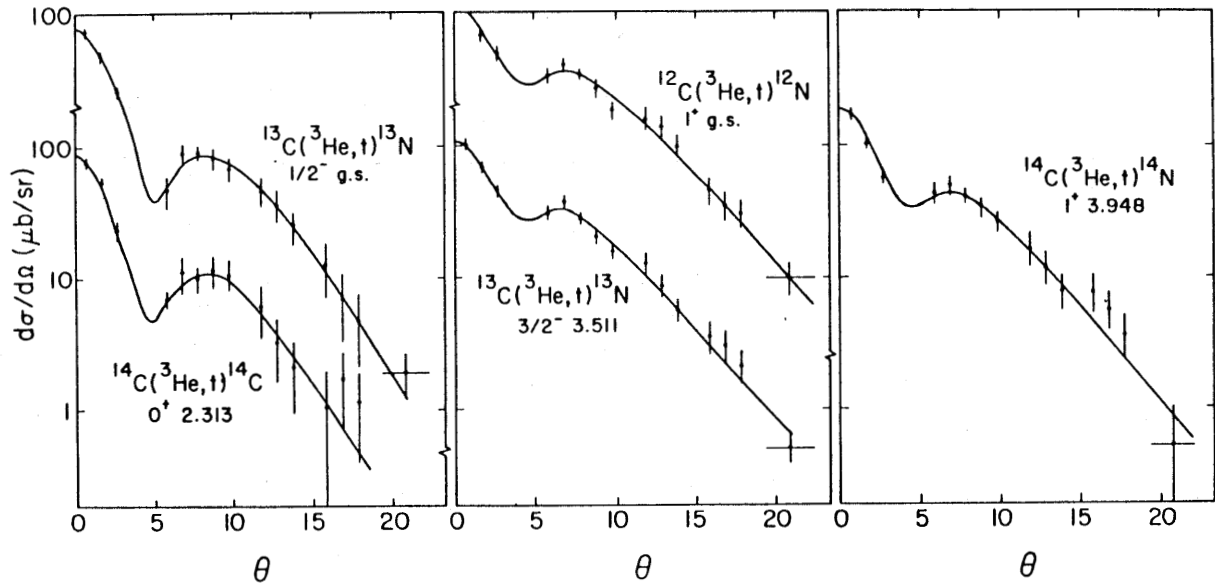


Figure 2. Angular distributions for Fermi (left) and Gamow-Teller (middle and right) transitions to low-excited states in  $^{12,13,14}\text{N}$ . The lines are drawn to guide the eye.