

CHARGE EXCHANGE

$D_{NN}(0^\circ)$ FOR THE $^{17,18}\text{O}(\vec{p}, \vec{n})^{17,18}\text{F}$ REACTIONS AT $E_{\vec{p}} = 120$ MeV

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For the $^{17}\text{O}(\vec{p}, \vec{n})^{17}\text{F}(\text{g.s.})$ transition, i.e., an odd mass $T = 1/2$ mirror transition, both GT and F amplitudes contribute to the cross section. The relative weighting of GT and Fermi cross sections in 0° (p,n) spectra can be parameterized as a function of bombarding energy in the 50 – 200 MeV region as

$$\frac{\sigma_{GT}/B(GT)}{\sigma_F/B(F)} = \left(\frac{E_p}{E_0} \right)^2 \equiv R^2 \quad (1)$$

where $E_0 = 55$ MeV is an empirically determined quantity. D_{NN} measurements for the $^{17}\text{O}(\vec{p}, \vec{n})^{17}\text{F}(\text{g.s.})$ were motivated by some unexpected fluctuations of the ratio R , especially in even-odd mass nuclei.^{1,2} Transverse polarization transfer measurements yield direct determinations of the relative weighting of the GT and F contributions. Therefore, D_{NN} measurements will allow an independent check on the aforementioned empirical relationship.

We are interested in using D_{NN} from the $^{18}\text{O}(\text{p,n})$ reaction to extract the cross section to the 0^- state at 1.08 MeV in ^{18}F . The IAS cross section to the 1.04 MeV state may include some contribution from the 0^- state, but we cannot hope to resolve the two. Since D_{NN}

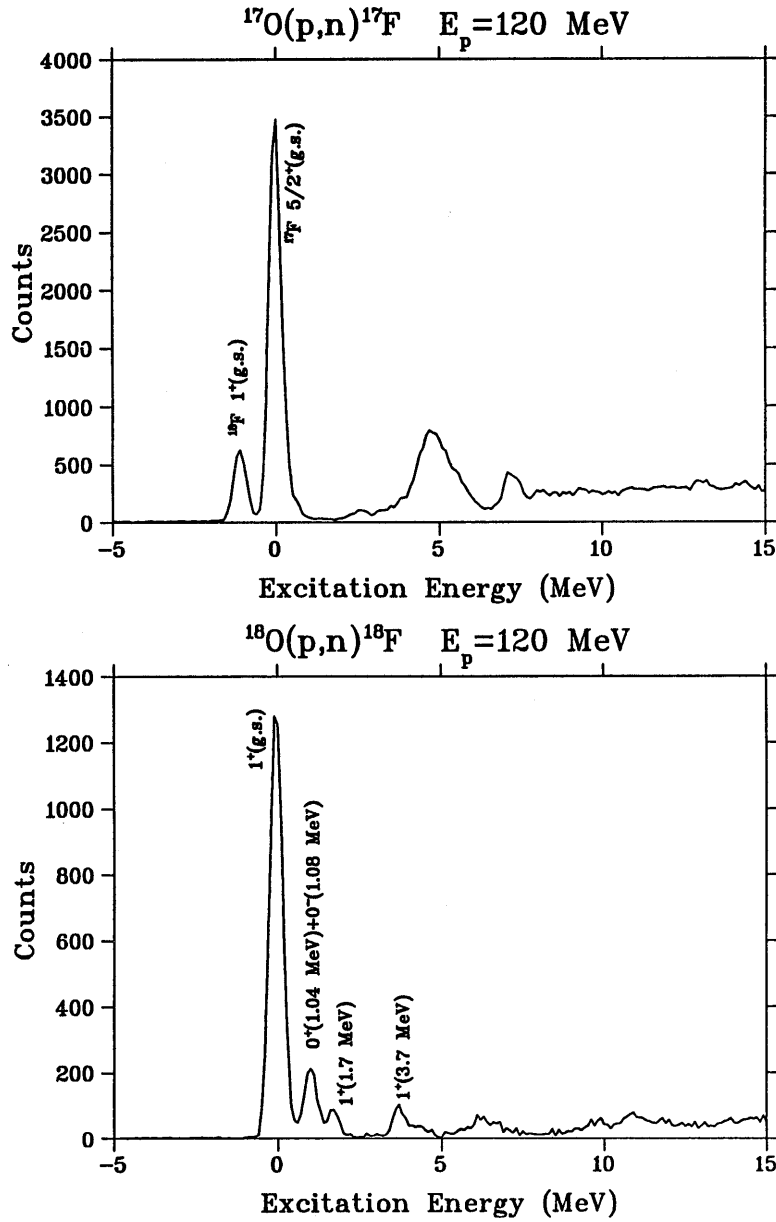


Figure 1. The 0° -spectra for the $^{17,18}\text{O}(\vec{p}, \vec{n})^{17,18}\text{F}$ reactions at $E_p = 120$ MeV are shown. The peak located to the left of the $^{17}\text{O}(\vec{p}, \vec{n})^{17}\text{F}(g.s.)$ is contamination from ^{18}O . The $^{17}\text{O}(\vec{p}, \vec{n})^{17}\text{F}(g.s.)$ peak also contains contaminations from low lying states in $^{18}\text{O}(p,n)$.

is +1 for the IAS and -1 for the 0^- state, a 0^- yield as small as 10% of the peak should be measurable (a 10% cross section is estimated from DWIA calculation). The ratio R is related to the parameters of the factorized DWIA expression for the (p,n) cross section through:

$$\frac{\hat{\sigma}_{GT}}{\hat{\sigma}_F} = \frac{\sigma_{GT}(q=0)}{\sigma_F(q=0)} \frac{B(F)}{B(GT)} \equiv R^2 = \left(\frac{E_p}{55} \right)^2 \quad (2)$$

$$\hat{\sigma}_\alpha = K(E_p) N_\alpha^D |J_\alpha|^2; \alpha = GT, F \quad (3)$$

where $K(E_p)$ and N_α^D are the kinematic factor and distortion factor, respectively, and J_α is the GT or F effective interaction potential integrated over the nuclear volume.

The experiment was performed at the IUCF Beam Swinger facility using a pressurized gas cell filled with oxygen. The $^{17}\text{O}_2$ target was a mixture of ^{16}O (17.7 atom%), ^{17}O (76.1 atom%), and ^{18}O (6.2 atom%); the $^{18}\text{O}_2$ target was 99% ^{18}O with a trace of ^{16}O . The gas cell was a welded stainless steel box with 25- μm thick Havar entrance and exit foils clamped between window frames which were also welded. At room temperature, 7 atm of pressure was easily obtained in the cell, giving about 50 mg/cm² of oxygen target thickness. An energy resolution of about 500 keV was achieved.

Preliminary excitation energy spectra for the $^{17,18}\text{O}(\vec{p}, \vec{n})^{17,18}\text{F}$ reactions are shown in Fig. 1. A preliminary D_{NN} value of -0.099 ± 0.004 was obtained for the $^{17}\text{O}(\vec{p}, \vec{n})^{17}\text{F}(\text{g.s.})$, where the error is statistical only. This preliminary result differs slightly from the expected value of -0.112. Data analysis is in progress to obtain D_{NN} for the peak that encompasses the 1.04 and 1.08 MeV states.

1 J. Wagner, Ph.D. thesis, Ohio State University (1989).

2 Y Wang, *et al.*, IUCF Sci. and Tech. Rep., May 1988 - June 1989, p. 48.