## THE 14N(p,p') REACTION AT 159.4 MeV

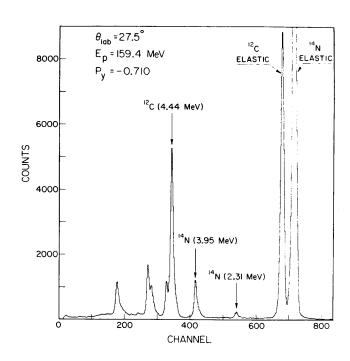
T.N. Taddeucci and J.R. Rapaport Ohio University, Athens, Ohio 45701

## J.R. Comfort University of Pittsburgh, Pittsburgh, Pennsylvania 15260

C.C. Foster
Indiana University Cyclotron Facility, Bloomington, Indiana 47405

The QDDM magnetic spectrograph and a 159.4 MeV polarized proton beam have been used at the IUCF to measure differential cross sections and analyzing powers for the  $^{14}N(p,p')^{14}N$  (g.s., 2.31-MeV and 3.95-MeV) transitions. The data span the angular range 7.5° (12.5°)  $\leq \theta_{1ab} \leq 45$ ° for the elastic (inelastic) transitions in 2.5° steps. Absolute normalization of the differential cross sections has been determined from the measured thickness and known composition of the melamine (C3H6N6) target. During the same run a graphite target was used to obtain data for the  $^{12}C(p,p')^{12}C$  (g.s. and 4.44-MeV) transitions. The graphite-target data allow subtraction of the 12C contribution to the melamine spectra (a problem only for the forward-angle elastic peak) and provide a check on the cross-section normalizations. The 12C cross sections obtained using the nominal graphite and melamine target thicknesses agree to within 5%. A typical spectrum obtained at  $\theta_{lab} = 27.5^{\circ}$  is shown in Fig. 1.

The differential cross section angular distributions for the transitions to the 2.31-MeV and 3.95-MeV states in  $^{14}N$  (Fig. 2) are similar to those previously measured at 122 MeV.1) Analyzing powers for these transitions are shown in Fig. 3. For the 2.31-MeV state,  $|A_y(\theta)|<0.3$  and the shape resembles that for the 15.11-MeV transition in  $^{12}C$  for momentum transfer q > 200 MeV/c.  $A_y(\theta)$  for the 3.95-MeV state is quite similar to that of the 4.44-MeV transition in  $^{12}C.2)$ 



<u>Figure 1</u>. Spin-down proton spectrum obtained at  $\theta_{1ab}$  = 27.5° with a melamine target.

The 2.31-MeV transition involves spin-parity and isospin transfers of  $\Delta J^{\pi}=1^+$ ,  $\Delta T=1$ . It is related through isospin invariance to the  $^{14}C(p,n)^{14}N$  (g.s.) and  $^{14}N(p,n)^{14}O(g.s)$  reactions. A preliminary angular distribution for the  $^{14}C(p,n)^{14}N(g.s)$  reaction is shown in Fig. 2. The (p,n) data are discussed in more detail elsewhere in this report. Time-reversal and isospin invariance arguments predict that the  $^{14}C(p,n)$  and  $^{14}N(p,p^*)$  cross sections should be related by  $^3)$ 

$$\frac{d\sigma}{d\Omega}$$
 (p,n) = 6  $\frac{d\sigma}{d\Omega}$  (p,p')

The absolute normalization at the (p,n) data in Fig. 2 is at present uncertain and has been adjusted to reflect the above relationship. Note that there is excellent agreement in shape between the angular distributions.

The (p,p') and (p,n) transitions between the 14Nground state and the  $0^+$ , T=1  $^{14}$ C(g.s.),  $^{14}$ N(2.31-MeV), and  $^{14}0(g.s)$  levels have been used as tests of tensor-force strength in the microscopic effective

○ - 3.95 MeV -2.31 MeV <sup>14</sup>C (p, n) <sup>14</sup>N (g.s.) IO\_X 10-1  $d\sigma/d\Omega_{c.m.}$  (mb/sr) 0 0 ′IŌ<sup>2</sup> 102 0

Figure 2. Differential cross sections for the transitions to the 2.31-MeV and 3.95-MeV states in  $^{14}\,\mathrm{N}.$ The  $^{14}C(p,n)$   $^{14}N(g.s.)$  points are from a preliminary analysis of data obtained by Goodman et al. (see contribution, this report).

30

 $\theta_{\rm c.m.}$ 

40

50

60

10

20

interaction. 1, 3, 4, 5, 6) The L=0 central-force contribution to the direct transition amplitudes for these reactions is suppressed in a manner analogous to the well-known retardation of the 14C Gamow-Teller  $\beta$ -decay rate.<sup>4)</sup> Microscopic distorted-wave calculations have met with mixed success in fitting differential cross sections for these transitions, with the best fits being obtained for  $E_p < 40 \text{ MeV.}^{3,5}$ ) DWIA calculations for comparison with the present data are under way.

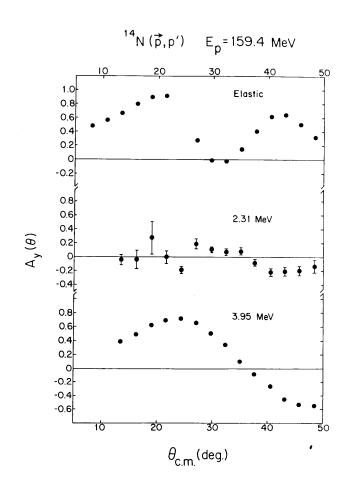


Figure 3. Analyzing powers for the transitions to the ground state and first two excited states in 14 N.

- J.R. Comfort, S. M. Austin, P.T. Debevec, G.L. Moake, R.W. Finlay, and W.G. Love, Phys. Rev. C21, 2147 (1980).
- J.R. Comfort, C.C. Foster, C.D. Goodman, D.W. Miller, G.L. Moake, P. Schwandt, J.R. Rapaport, R.E. Segal, Contribution to the Fifth International Symposium on Polarization Phenomena in Nuclear Physics, Santa Fe, 1980.
- T.N. Taddeucci, R.R. Doering, L.C. Dennis, S.M. Austin, A. Galonsky, W.M. Sterrenburg, preprint and Bull. Am. Phys. Soc. 25, 728 (1980).
- 4) C. Wong, J.D. Anderson, V.A. Madsen, F.A. Schmittroth, and M.J. Stomp, Phys. Rev. C3, 1904 (1971).
- 5) S.H. Fox and S.M. Austin, Phys. Rev. <u>C24</u>, 1133 (1980).
- W.G. Love, L.J. Parish, A. Richter, Phys. Lett. 31B, 167 (1970).

ANALYZING POWERS FOR THE PROTON EXCITATION OF HIGH-SPIN STATES IN <sup>28</sup>S1:
A NEW LOOK AT THE EFFECTIVE INTERACTION

A.D. Bacher, G.T. Emery, W.P. Jones, D.W. Miller, C. Olmer, P. Schwandt Indiana University Cyclotron Facility, Bloomington, Indiana 47405

S. Yen, R.J. Sobie and <u>T.E. Drake</u>
Department of Physics, University of Toronto, Toronto, Canada M5S 1A7

W.G. Love
Department of Physics and Astronomy, University of Georgia,
Athens, Georgia 30602

F. Petrovich
Department of Physics, The Florida State University,
Tallahassee, Florida 32306

One of the outstanding new features of inelastic proton scattering at intermediate energies has been the observation of high-spin states of simple structure, that are excited primarily through the non-central components of the effective nucleon-nucleon interaction.  $^{1-6}$ ) In the present work, analyzing powers  $A_y(\theta)$  have been measured for 135 MeV (p,p') excitation of the 5-, T=0 (9.70 MeV), 6-, T=0 (11.58 MeV), and 6-, T=1 (14.35 MeV) states in  $^{28}$ Si. Differential cross sections for transitions of this nature have provided definitive information on the strength of the high-momentum components of the non-central parts of the effective interaction,  $^{4-6}$ ) and generally confirm the validity of the impulse approximation.  $^{7}$ 

For the three states of interest, the differential cross sections, analyzing powers, and their products are shown from top to bottom in Figs. 1 and 2. While the shapes of the cross-section angular distributions are quite similar, the analyzing-power angular distributions are distinctly different and provide a definite signature of the spin and isospin transfer for each transition. The curves in Figs. 1 and 2 are the result of distorted-wave impulse approximation (DWIA) calculations. These calculations used the complex central and real spin-orbit components of the two-nucleon t-matrix interaction<sup>8</sup>) supplemented by a real tensor interaction. The imaginary parts of t<sup>LS</sup> and t<sup>T</sup>, which have been neglected, are small. Optical-model parameters were taken from Schwandt